

Experimental Confirmation of n —scroll Hyperchaotic Attractors

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Abstract—A systematic circuit design approach is proposed for experimental verification of hyperchaotic 2, 3, 4—scroll attractors from a generalized Matsumoto-Chua-Kobayashi (MCK) circuit. The recursive formulas for system parameters are rigorously derived for improving the hardware implementation.

I. INTRODUCTION

Hyperchaos was first observed from a real physical system by Matsumoto, Chua and Kobayashi in [1]. Then, Yalcin et al. [2] introduced some hyperchaotic n —double-scroll chaotic attractors by adding breakpoints in the piecewise-linear (PWL) characteristic of the MCK circuit and confirmed the hyperchaotic 4— and 6—scroll attractors by computer simulations. Yu et al. [3] proposed hyperchaotic n —scroll attractors and realized hyperchaotic 3 ~ 10—scroll attractors by computer simulations. Itoh et al. [4] investigated the impulsive synchronization of a hyperchaotic double-scroll attractor and its application to spread-spectrum communication systems. It has been known that it is generally difficult to implement multi-scroll chaotic and hyperchaotic attractors by a physical electronic circuit. Yalcin et al. [5] experimentally confirmed 3— and 5—scroll chaotic attractors in a generalized Chua's circuit, while Zhong et al. [6] proposed a systematical circuitry design method for physically implementing up to as many as ten scrolls visible on the oscilloscope. Han et al. [7] constructed a double-hysteresis building block to physically realize a 9—scroll chaotic attractor. There are some other approaches reported in the literature for the design and circuit implementation of multi-scroll chaotic attractors [8-14]. It is generally quite difficult to physically build a nonlinear resistor having an appropriate characteristic with many segments. In this effort, Lü et al. [13] designed a novel circuit diagram to physically verify the multi-directional multi-scroll chaotic attractors. The main obstacle is that the device must have a very wide dynamic range [3,6], however physical conditions always limit or even prohibit such circuit realization [6]. Recently, Lü and Chen [14] reviewed the main advances of multi-scroll chaos generation.

In this paper, we describe the design of a novel block circuit diagram to experimentally confirm hyperchaotic n —scroll attractors. This is the first time in the literature to report

an experimental verification of hyperchaotic 3— and 4—scroll attractors. Moreover, the derived recursive formulas for system parameters provide a theoretical basis for physical realization of hyperchaotic attractors with a large number of scrolls.

The rest of the paper is organized as follows. In Section II, a general MCK circuit is briefly described. Then, a novel block circuit diagram is designed for hardware implementation of hyperchaotic 2, 3, 4—scroll attractors, and its dynamic equation is rigorously derived in Section III. Conclusions are finally drawn in Section IV.

II. A GENERALIZED MCK CIRCUIT

The dimensionless state equation of the hyperchaotic MCK circuit is described by [1]

$$\begin{cases} \frac{dx}{d\tau} &= \alpha[g(y-x) - z] \\ \frac{dy}{d\tau} &= \beta[-g(y-x) - w] \\ \frac{dz}{d\tau} &= \gamma_0(x+z) \\ \frac{dw}{d\tau} &= \gamma y, \end{cases} \quad (1)$$

where $g(y-x) = m_1(y-x) + 0.5(m_0 - m_1)[|y-x+1| - |y-x-1|]$. When $\alpha = 2$, $\beta = 20$, $\gamma_0 = 1$, $\gamma = 1.5$, $m_0 = -0.2$, $m_1 = 3$, system (1) has a hyperchaotic double-scroll attractor with Lyapunov exponents $\lambda_1 = 0.24$, $\lambda_2 = 0.06$, $\lambda_3 = 0$, $\lambda_4 = -53.8$.

To generate hyperchaotic n —scroll attractors from (1), we first generalize the characteristic function $g(y-x)$, given in [3], as follows:

$$g(y-x) = m_{N-1}(y-x) + 0.5 \sum_{i=1}^{N-1} (m_{i-1} - m_i)(|y-x+x_i| - |y-x-x_i|). \quad (2)$$

The recursive formulas of positive switching points x_i ($i =$

2, 3, ..., N - 1) can be easily deduced as follows:

$$\begin{cases} x_2 = \frac{(1+k_1) \sum_{i=1}^1 (m_i - m_{i-1}) x_i}{m_1 - 1} - k_1 x_1 \\ x_3 = \frac{(1+k_2) \sum_{i=1}^2 (m_i - m_{i-1}) x_i}{m_2 - 1} - k_2 x_2 \\ \vdots \\ x_{N-1} = \frac{(1+k_{N-2}) \sum_{i=1}^{N-2} (m_i - m_{i-1}) x_i}{m_{N-2} - 1} - k_{N-2} x_{N-2}, \end{cases} \quad (3)$$

where $m_i (0 \leq i \leq N-1)$ are the slopes of the segments and radials in various PWL regions, and $k_i = \frac{x_{i+1} - x_i^E}{x_i^E - x_i} (1 \leq i \leq N-2)$, in which $x_i^E (1 \leq i \leq N-2)$ are the positive equilibrium points of $g(x)$.

To control the hyperchaotic signal into the region of the operational amplifier, we may assume that $x_1 < 1$. Here, we suppose that $x_1 = 0.5$. From (3), we determine the system parameters as follows: (i) when $N = 2$, $m_0 = -0.2$, $m_1 = 3$, system (1) with (2) has a hyperchaotic double-scroll attractor; (ii) when $N = 3$, $m_0 = 3$, $m_1 = -0.8$, $m_2 = 3$, $x_2 = 1.8333$, system (1) with (2) has a hyperchaotic 3-scroll attractor; (iii) when $N = 4$, $m_0 = m_2 = -0.7$, $m_1 = m_3 = 2.9$, $x_2 = 1.5289$, $x_3 = 3.0239$, system (1) with (2) has a hyperchaotic 4-scroll attractor.

III. CIRCUIT DESIGN AND IMPLEMENTATION

In this section, a circuit diagram is constructed to experimentally verify the hyperchaotic 2, 3, 4-scroll attractors. Also, the dynamic equation is rigorously derived from the circuit diagram shown in Fig. 1.

A. Circuit diagram and its dynamic equation

Fig. 1 shows the circuit diagram, where N_1 is the generator of the negative resistor $-R$, and N_R is the multi-PWL function generator satisfying $I_N = f(v_{C2} - v_{C1})$. All operational amplifiers are selected as Type TL082. The voltage of the electric source is $E = 15V$. Thus, the saturating voltages of the operation amplifiers are $E_{sat} = 14.3V$.

According to Fig. 1, the circuit equation is derived as follows:

$$\begin{cases} C_1 \frac{dv_{C1}}{dt} = f(v_{C2} - v_{C1}) - i_{L1} \\ C_2 \frac{dv_{C2}}{dt} = -f(v_{C2} - v_{C1}) - i_{L2} \\ L_1 \frac{di_{L1}}{dt} = v_{C1} + R i_{L1} \\ L_2 \frac{di_{L2}}{dt} = v_{C2}, \end{cases} \quad (4)$$

where $f(v_{C2} - v_{C1}) = G_{N-1}(v_{C2} - v_{C1}) + 0.5 \sum_{i=1}^{N-1} (G_{i-1} - G_i)(|v_{C2} - v_{C1} + E_i| - |v_{C2} - v_{C1} - E_i|)$ is a piecewise-linear characteristic function.

Comparing systems (1) with (4), we get the transformation relationship of parameters as follows:

$$\begin{cases} \tau_0 = 2RC_1, \tau = \frac{t}{\tau_0}, \alpha = 2, \beta = \frac{2C_1}{C_2} = 20 \\ \gamma_0 = \frac{2R^2C_1}{L_1} = 1, \gamma = \frac{2R^2C_1}{L_2} = 1.5 \\ x = \frac{v_{C1}}{V_{BP}}, y = \frac{v_{C2}}{V_{BP}}, z = \frac{Ri_{L1}}{V_{BP}}, w = \frac{Ri_{L2}}{V_{BP}} \\ V_{BP} = 1V, G_i = m_i G (i = 0, 1, 2, \dots), G = \frac{1}{R} \\ g(y - x) = Rf(v_{C2} - v_{C1}), \end{cases} \quad (5)$$

where $V_{BP} = 1V$, $\frac{1}{\tau_0} = \frac{1}{2RC_1}$ is the time-scale transformation factor.

From (5), we have the parameters: $L_1 = 9mH$, $L_2 = 6mH$, $C_1 = 50nF$, $C_2 = 5nF$, $R = 300\Omega$. Then, we can get the theoretical values of the resistors based on the parameters given in Section II as follows:

(1) For hyperchaotic 2-scroll attractor:

$$\begin{cases} G_0 = \frac{m_0}{R} = -0.67mS, G_1 = \frac{m_1}{R} = 10mS, \\ E_1 = x_1 V_{BP}, r_1 = \frac{R_{12}}{R_{11}} = G_1 R_2 - 1 = 1.00, \\ r_2 = \frac{R_{22}}{R_{21}} = \frac{E_{sat}}{E_1} = 28.6, \\ r_3 = \frac{R_{32}}{R_{31}} = \frac{r_2}{R_2(G_1 - G_0)} - 1 = 12.4. \end{cases} \quad (6)$$

(2) For hyperchaotic 3-scroll attractor:

$$\begin{cases} G_0 = \frac{m_0}{R} = 10mS, G_1 = \frac{m_1}{R} = -2.7mS, \\ G_2 = \frac{m_2}{R} = 10mS, E_i = x_i V_{BP} (i = 1, 2), \\ r_1 = \frac{R_{12}}{R_{11}} = G_2 R_2 - 1 = 1.00, \\ r_2 = \frac{R_{22}}{R_{21}} = \frac{E_{sat}}{E_2} = 7.80, \\ r_3 = \frac{R_{32}}{R_{31}} = \frac{r_2}{R_2(G_2 - G_1)} - 1 = 2.08, \\ r_4 = \frac{R_{42}}{R_{41}} = \frac{E_{sat}}{E_1} - 1 = 27.60, \\ r_5 = \frac{R_{52}}{R_{51}} = -\frac{1+r_4}{R_2(G_1 - G_0)} - 1 = 10.29. \end{cases} \quad (7)$$

(3) For hyperchaotic 4-scroll attractor:

$$\begin{cases} G_0 = \frac{m_0}{R} = -2.3mS, G_1 = \frac{m_1}{R} = 9.7mS, \\ G_2 = \frac{m_2}{R} = -2.3mS, G_3 = \frac{m_3}{R} = 9.7mS, \\ E_i = x_i V_{BP} (i = 1, 2, 3), r_1 = \frac{R_{12}}{R_{11}} = G_3 R_2 - 1 = 0.93, \\ r_2 = \frac{R_{22}}{R_{21}} = \frac{E_{sat}}{E_3} = 4.73, \\ r_3 = \frac{R_{32}}{R_{31}} = \frac{r_2}{R_2(G_3 - G_2)} - 1 = 0.97, \\ r_4 = \frac{R_{42}}{R_{41}} = \frac{E_{sat}}{E_2} - 1 = 8.35, \\ r_5 = \frac{R_{52}}{R_{51}} = -\frac{1+r_4}{R_2(G_2 - G_1)} - 1 = 2.90, \\ r_6 = \frac{R_{62}}{R_{61}} = \frac{E_{sat}}{E_1} = 28.6, \\ r_7 = \frac{R_{72}}{R_{71}} = \frac{r_6}{R_2(G_1 - G_0)} - 1 = 10.90. \end{cases} \quad (8)$$

B. Experimental observations

TABLE I
THE RATIOS OF THE RESISTORS $r_n = \frac{R_{n2}}{R_{n1}} (1 \leq n \leq 7)$
AND THE NUMBER OF THE SCROLLS N

r_1	r_2	r_3	r_4	r_5	r_6	r_7	N
1.00	28.60	12.40					2
1.00	7.80	2.08	27.60	10.29			3
0.93	4.73	0.97	8.35	2.90	28.60	10.90	4

TABLE II
THE RESISTORS $R_{n2} = r_n R_{n1} (1 \leq n \leq 7)$ AND THE
NUMBER OF THE SCROLLS N

R_{12}	R_{22}	R_{32}	R_{42}	R_{52}	R_{62}	R_{72}	N
10k	286k	12.4k					2
10k	78k	2.08k	276k	10.29k			3
9.3k	47.3k	0.97k	83.5k	2.90k	286k	10.9k	4

Let $R_1 = 100k\Omega$, $R_2 = 0.2k\Omega$, $R_{31} = R_{51} = R_{71} = 1k\Omega$, $R_{11} = R_{21} = R_{41} = R_{61} = 10k\Omega$. By comparing Fig. 1 with system (1) under (2), we can calculate the resistors $R_{n2} (1 \leq n \leq 7)$ as shown in Tables I and II. As seen from Fig. 1, when K_1, K_2 are switched on and K_3, K_4 are

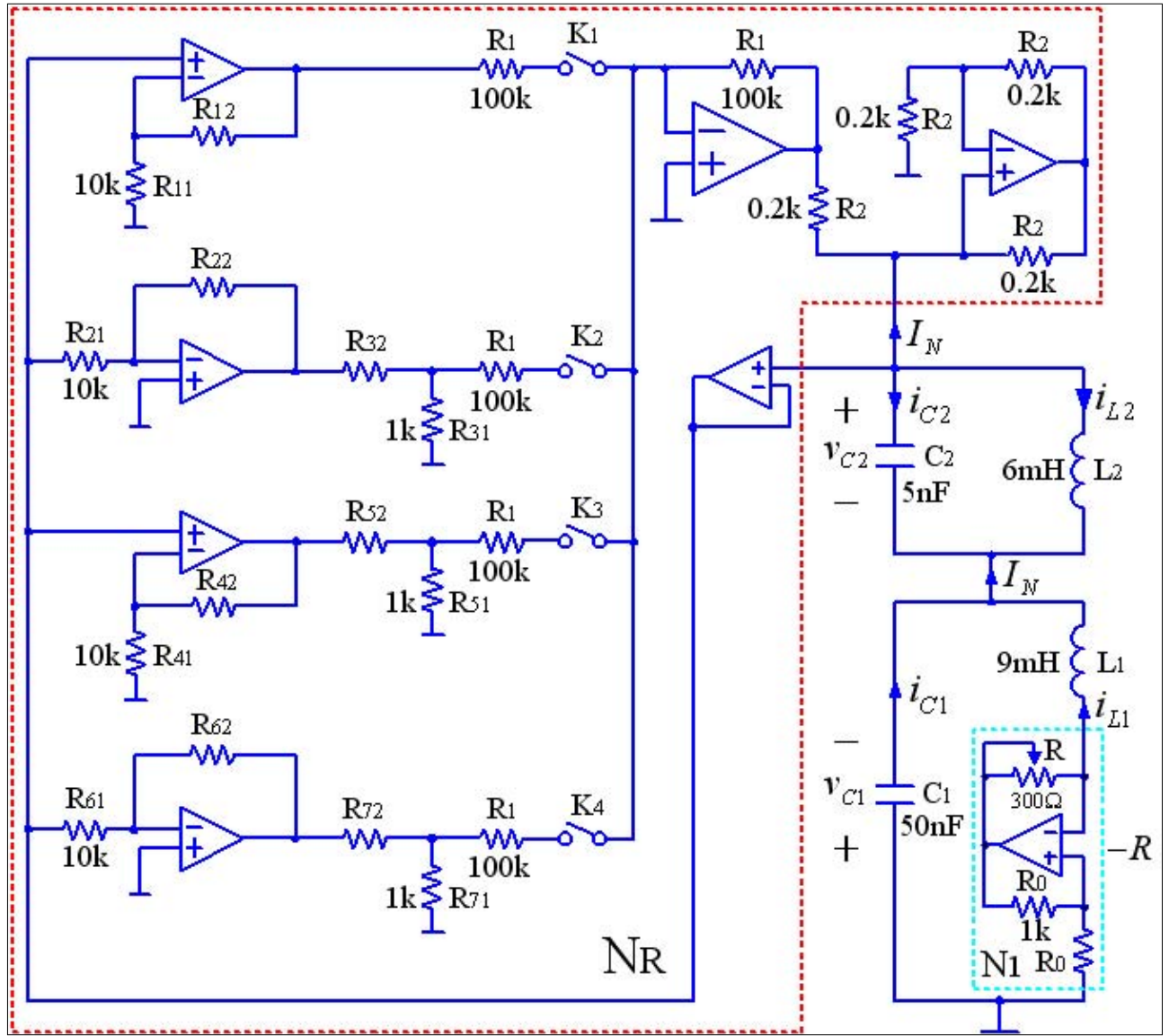


Fig. 1. Circuit diagram for generating hyperchaotic n -scroll attractors.

switched off, the circuit diagram can create a hyperchaotic double-scroll attractor; when K_1 , K_2 , K_3 are switched on and K_4 is switched off, the circuit diagram can generate a hyperchaotic 3-scroll attractor, as shown in Fig. 2 (a); when K_1 , K_2 , K_3 , K_4 are switched on, the circuit diagram can create a hyperchaotic 4-scroll attractor, as shown in Fig. 2 (b).

IV. CONCLUSIONS

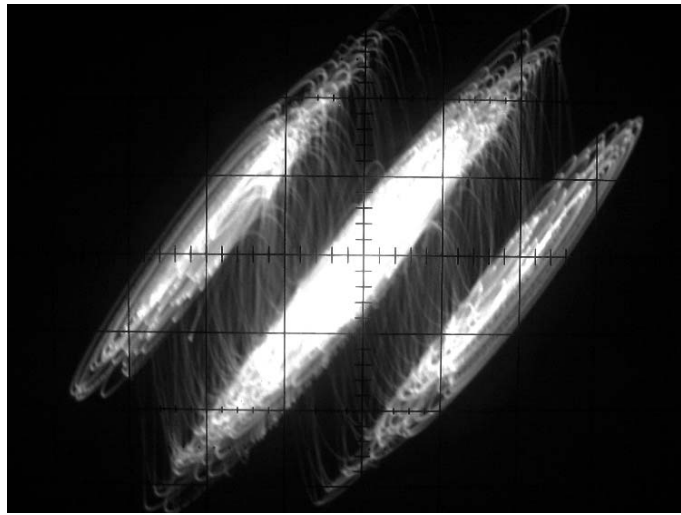
This brief paper has proposed a novel block circuit diagram for hardware implementation of hyperchaotic 2, 3, 4-scroll attractors in a generalized MCK circuit. In addition, the derived recursive formulas for system parameters provide a theoretical basis for physical realization of the hyperchaotic attractors with a large number of scrolls.

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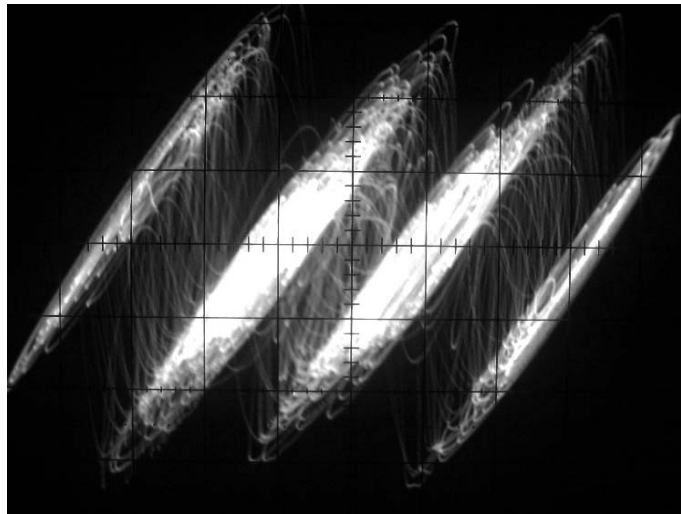
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(a) 3—scroll



(b) 4—scroll

Fig. 2. Experimental observations of hyperchaotic n —scroll attractors.

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