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The Stabilization of a One-Dimensional Wave Equation by Boundary Feedback With Noncollocated Observation

Bao-Zhu Guo and Cheng-Zhong Xu

Abstract—This note addresses the stabilization of a one-dimensional wave equation with control at one end and noncollocated observation at another end. A simple exponentially convergent observer is constructed. The dynamical stabilizing boundary output feedback is designed via the observed state. While the closed-loop system is nondissipative, we show its exponential stability using Riesz basis approach.

Index Terms—Exact controllability, exponential stability, feedback control, observer, wave equation.

I. INTRODUCTION

It has been known by engineers for a long time that a partial differential equation describing a mechanical system, like a flexible structure in which the power flow into the system is the scalar product $\langle u, y \rangle$ (e.g., when u is force and y is velocity), leads to a positive-real system if actuators and sensors are designed in a "collocated" fashion. Collocated systems always relate with "passivity" which was introduced in connection with circuit theory in 1950s [6]. By "passivity," we mean that the increase of energy stored in the system does not exceed the energy that enters from the external world. For such a system, the transfer function is positive real and negative output feedback produces a dissipative system, which is stable in the sense of Lyapunov.

However, collocated control design is not always feasible in practice and its performance is not always good enough [2]. Actually, the noncollocated control has been widely used in engineering systems control, see [1], [12], [10], [15], [18], [20], [21], [23], [25], and [27]. It is well-known that the noncollocated systems are usually not minimum-phase. As a result, a small increment of feedback controller gains can easily make the closed-loop system unstable [18]. Hence, the control design for the stabilization of noncollocated systems is much harder than collocated ones. Compared with the huge works on the stabilization of collocated PDEs in literature, the study for noncollocated

PDEs (systems described by partial differential equations) control is still fairly scarce. Very recently, the observed state feedbacks are designed through backstepping observers in [19] to stabilize a class of one-dimensional parabolic PDEs. The abstract observers design for a class of well-posed regular infinite-dimensional systems can be found in [5] but the stabilization is not addressed.

In this note, a simple observer for a one-dimensional wave equation with boundary control and noncollocated observation is constructed. This observer design is closely related with the theory of well-posed linear infinite-dimensional systems that has been studied extensively in the last two decades [8]. This is because that we require output of the system and state of the observer to depend continuously on input. This seems very natural in design of the observers. The state space for our observer is larger than the state space of the original system. We prove that this observer is exponentially convergent. The boundary output feedback control is designed by state of the observer. This results in nondissipativity of the closed-loop system, which gives rise to difficulty of proof for the stability. Fortunately, using Riesz basis approach, we show that the closed-loop system is indeed exponentially stable. Another difficulty that should be pointed out is that the uncontrolled wave equation has infinite many of eigenvalues on the imaginary axis. This is different to the parabolic equation that has at most a finite number of eigenvalues on the closed right-half complex plane, for which the LQ approach can also be used to solve the problem [13]. Furthermore, our problem is also not the regular in the sense of [5] where the feedthrough operator is required to be zero.

The problem we are concerned with is the following one-dimensional wave equation with boundary control and noncollocated observation:

$$\begin{cases} w_{tt}(x, t) - w_{xx}(x, t) = 0, & 0 < x < 1, t > 0 \\ w(0, t) = 0, & t \geq 0 \\ w_x(1, t) = u(t), & t \geq 0 \\ y(t) = w_x(0, t), & t \geq 0 \end{cases} \quad (1.1)$$

where u is the control (input) and y is the observation (output). The state space is $H = H_L^1(0, 1) \times L^2(0, 1)$, $H_L^1(0, 1) = \{f \mid f \in H^1(0, 1), f(0) = 0\}$. H is equipped with the norm $\|(f, g)\|_H^2 = \int_0^1 [|f'(x)|^2 + |g(x)|^2] dx$ for any $(f, g) \in H$. The input (output) space is $U = \mathbb{C}$. A simple spectral analysis shows that for any direct PI output feedback $u = k_1 y + k_2 \int_0^t y(s) ds$ with reals k_1, k_2 , the closed loop system is always unstable.

Define the operator $\mathcal{A} : D(\mathcal{A}) \subset H \rightarrow H$ as following:

$$\begin{aligned} \mathcal{A}(f, g) &= (g, f'') \\ D(\mathcal{A}) &= \{(f, g) \in H \mid \mathcal{A}(f, g) \in H, f'(1) = 0\}. \end{aligned} \quad (1.2)$$

Then, (1.1) can be written as

$$\sum(\mathcal{A}, \mathcal{B}, \mathcal{C}) : \begin{cases} \frac{d}{dt} \begin{pmatrix} w \\ w_t \end{pmatrix} = \mathcal{A} \begin{pmatrix} w \\ w_t \end{pmatrix} + \mathcal{B}u(t) \\ y(t) = \mathcal{C} \begin{pmatrix} w \\ w_t \end{pmatrix} = w_x(0, t) \end{cases} \quad (1.3)$$

where

$$\mathcal{B} = \begin{pmatrix} 0 \\ \delta(x-1) \end{pmatrix}, \quad \mathcal{C} = (-\langle \delta'(x), \cdot \rangle, 0)$$

and $\delta(\cdot)$ denotes the Dirac function. Obviously, both \mathcal{B} and \mathcal{C} are unbounded operators.

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Theorem 2: For any $u(\cdot), w_x(0, \cdot) \in L^2_{loc}(0, \infty)$ and initial datum $(\hat{w}(\cdot, 0), \hat{w}_t(\cdot, 0)) \in \mathbf{H}$, there exists a unique solution $(\hat{w}, \hat{w}_t) \in C(0, \infty; \mathbf{H})$ to (2.1), and for all $T > 0$, there exists a $D_T > 0$ depending on T only such that

$$\|(\hat{w}(\cdot, T), \hat{w}_t(\cdot, T))\|_{\mathbf{H}}^2 \leq D_T \left\{ \|(\hat{w}(\cdot, 0), \hat{w}_t(\cdot, 0))\|_{\mathbf{H}}^2 + \int_0^T [|u(\tau)|^2 + |w_x(0, \tau)|^2] d\tau \right\}. \quad (2.6)$$

Proof: It suffices to show that \mathbf{B}^* is admissible for $e^{\mathbf{A}^*t}$ (see [8, Def. 2.1] and [24, Th. 6.9]). A simple computation shows that this is equivalent to saying that $\mathbf{B}^* \mathbf{A}^{*-1}$ is bounded and for any $T > 0$, there exists a $M_T > 0$ depending on T only such that the system of the following:

$$\begin{cases} \hat{w}_{tt}(x, t) - \hat{w}_{xx}(x, t) = 0, & 0 < x < 1, \quad t > 0 \\ \hat{w}_x(0, t) = \alpha \hat{w}_t(0, t) + \beta \hat{w}(0, t), & t \geq 0 \\ \hat{w}_x(1, t) = 0, & t \geq 0 \\ y_w(t) = (\hat{w}_t(1, t), -\hat{w}_t(0, t)), & t \geq 0 \end{cases} \quad (2.7)$$

satisfies (see also [16, eq. (2.3)])

$$\int_0^T [|\hat{w}_t(1, t)|^2 + |\hat{w}_t(0, t)|^2] dt \leq M_T E_m(0)$$

where

$$E_m(t) = \frac{1}{2} \int_0^1 [|\hat{w}_x(x, t)|^2 + |\hat{w}_t(x, t)|^2] dx + \frac{\beta}{2} |\hat{w}(0, t)|^2.$$

First, a simple computation shows that

$$\begin{aligned} \mathbf{A}^{*-1}(\phi, \psi) &= \left(-\int_1^x (x-\tau)\psi(\tau) d\tau + \int_0^1 (\beta^{-1} + \tau)\psi(\tau) d\tau - \frac{\alpha}{\beta}\phi(0), -\phi(x) \right) \quad \forall (\phi, \psi) \in \mathbf{H} \\ \mathbf{B}^* \mathbf{A}^{*-1}(\phi, \psi) &= (\phi(1), -\phi(0)) \quad \forall (\phi, \psi) \in \mathbf{H}. \end{aligned}$$

Hence $\mathbf{B}^* \mathbf{A}^{*-1}$ is bounded on \mathbf{H} . Second, we consider once again the real function case only. Now, differentiate $E_m(t)$ in t to give

$$\dot{E}_m(t) = -\alpha \hat{w}_t^2(0, t) \leq 0$$

and, hence, $E_m(T) \leq E_m(0)$ for any $T > 0$ and

$$\int_0^T \hat{w}_t^2(0, t) dt \leq \frac{1}{\alpha} E_m(0).$$

Next, let

$$\rho_0(t) = \int_0^1 x \hat{w}_x(x, t) \hat{w}_t(x, t) dx$$

Then, $|\rho_0(t)| \leq E_m(t)$. Since

$$\dot{\rho}_0(t) = \frac{1}{2} \hat{w}_t^2(1, t) - \frac{1}{2} \int_0^1 [\hat{w}_x^2(x, t) + \hat{w}_t^2(x, t)] dx$$

it follows that

$$\int_0^T \hat{w}_t^2(1, t) dt \leq (2+T)E_m(0) \quad \forall T > 0.$$

The result is proved. \square

To end this section, we show that the observer is exponentially convergent in \mathbf{H} . Actually, let

$$z(x, t) = \hat{w}(1-x, t) - w(1-x, t).$$

Then z satisfies

$$\begin{cases} z_{tt}(x, t) - z_{xx}(x, t) = 0, & 0 < x < 1, \quad t > 0 \\ z_x(0, t) = 0, & t \geq 0 \\ z_x(1, t) = -\alpha z_t(1, t) - \beta z(1, t), & t \geq 0. \end{cases} \quad (2.8)$$

It is well-known that the system (2.8) associates with a C_0 -semi-group solution in \mathbf{H} and is exponentially stable in the sense of

$$E_\varepsilon(t) \leq M e^{-\omega t} E_\varepsilon(0)$$

for some $M, \omega > 0$, where

$$E_\varepsilon(t) = \frac{1}{2} \int_0^1 [|z_x(x, t)|^2 + |z_t(x, t)|^2] dx + \frac{\beta}{2} |z(1, t)|^2.$$

III. OUTPUT FEEDBACK STABILIZATION

In this section, we design the output feedback control law

$$u(t) = -\alpha \hat{w}_t(1, t)$$

for the system (1.1) and (2.1). The closed-loop system becomes

$$\begin{cases} \hat{w}_{tt}(x, t) - \hat{w}_{xx}(x, t) = 0, & 0 < x < 1, \quad t > 0 \\ \hat{w}_x(0, t) = \alpha \hat{w}_t(0, t) + \beta \hat{w}(0, t) + w_x(0, t), & t \geq 0 \\ \hat{w}_x(1, t) = -\alpha \hat{w}_t(1, t), & t \geq 0 \\ w_{tt}(x, t) - w_{xx}(x, t) = 0, & 0 < x < 1, \quad t > 0 \\ w(0, t) = 0, & t \geq 0 \\ w_x(1, t) = -\alpha \hat{w}_t(1, t), & t \geq 0. \end{cases} \quad (3.1)$$

Consider (3.1) in the space $X = \mathbf{H} \times H$ with norm induced by the inner product

$$\|(f, g, \phi, \psi)\|^2 = \int_0^1 [|f'(x)|^2 + |g(x)|^2 + |\phi'(x)|^2 + |\psi(x)|^2] dx + \beta |f(0)|^2 \quad \forall (f, g, \phi, \psi) \in X.$$

The system operator $\mathbb{A}(\subset X) \rightarrow X$ for (3.1) is defined by

$$\begin{cases} \mathbb{A}(f, g, \phi, \psi) = (g, f'', \psi, \phi'') & \forall (f, g, \phi, \psi) \in D(\mathbb{A}) \\ D(\mathbb{A}) = \{(f, g, \phi, \psi) \in X \mid \mathbb{A}(f, g, \phi, \psi) \in X \\ f'(0) = \alpha g(0) + \beta f(0) + \phi'(0) \\ \phi(0) = 0, f'(1) = \phi'(1) = -\alpha g(1)\}. \end{cases} \quad (3.2)$$

A simple computation shows that \mathbb{A} is not dissipative. So we use Riesz basis approach in the sequel to treat the C_0 -semigroup generation and stability property of \mathbb{A} .

Lemma 1: \mathbb{A}^{-1} is compact on X and hence $\sigma(\mathbb{A})$, the spectrum of \mathbb{A} , consists of isolated eigenvalues only.

Proof: For any $(p_1, q_1, p_2, q_2) \in X$, solve $\mathbb{A}(f, g, \phi, \psi) = (p_1, q_1, p_2, q_2)$ to obtain

$$\begin{cases} g(x) = p_1(x), & \psi(x) = p_2(x) \\ f(x) = \int_1^x (x - \tau) q_1(\tau) d\tau - \int_0^1 \tau q_1(\tau) d\tau - \alpha p_1(1)x + f(0) \\ f(0) = \frac{1}{\beta} \int_0^1 [q_2(\tau) - q_1(\tau)] d\tau - \frac{\alpha}{\beta} p_1(0) \\ \phi(x) = \int_1^x (x - \tau) q_2(\tau) d\tau - \int_0^1 \tau q_2(\tau) d\tau - \alpha p_1(1)x. \end{cases}$$

So, \mathbb{A}^{-1} is defined on whole space X and \mathbb{A}^{-1} maps X into a subset of space $(H^2(0, 1) \times H^1(0, 1))^2$, which is compact in X . By the Sobolev embedding theorem [14], \mathbb{A}^{-1} is compact on X , proving the required result. \square

Lemma 2: There are two families of eigenvalues of \mathbb{A} which can be expressed asymptotically as

$$\begin{cases} \lambda_{1n} = \frac{1}{2} \log \left| \frac{1-\alpha}{1+\alpha} \right| + n_\alpha \pi i, & \text{Re } \lambda_{1n} < 0, \quad n \in \mathbb{Z} \\ \lambda_{2n} = \frac{1}{2} \log \left| \frac{1-\alpha}{1+\alpha} \right| + n^\alpha \pi i + o\left(\frac{1}{|n|}\right), & |n| \rightarrow \infty \\ \lambda_{2n} < 0 \end{cases} \quad (3.3)$$

where

$$\begin{cases} n_\alpha = \begin{cases} n - 1/2, & 0 < \alpha < 1 \\ n, & \alpha > 1 \end{cases} \\ n^\alpha = \begin{cases} n, & 0 < \alpha < 1 \\ n - 1/2, & \alpha > 1. \end{cases} \end{cases} \quad (3.4)$$

The corresponding eigenfunctions for $\alpha \in (0, 1)$ are (there are parallel results for $\alpha > 1$)

$$\begin{cases} W_{1n}(x) = (F_n(x), F_n(x)), & n \in \mathbb{Z} \\ W_{2n}(x) = (\tilde{F}_n(x) + G_n(x), \tilde{F}_n(x)) \end{cases} \quad (3.5)$$

where

$$\begin{cases} G_n(x) = c_{2n} (\lambda_{2n}^{-1} p_{2n}(1-x), p_{2n}(1-x)) \\ c_{2n} = -\frac{1}{2\alpha} e^{\lambda_{2n}} \left[1 + \alpha + (1-\alpha) \frac{(1+\alpha)\lambda_{2n} + \beta}{(1-\alpha)\lambda_{2n} - \beta} \right] \\ = -\frac{\sqrt{1-\alpha^2}}{\alpha} (-1)^n + o\left(\frac{1}{|n|}\right) \end{cases} \quad (3.6)$$

$$\begin{cases} p_{2n}(x) = e^{\lambda_{2n}x} + e^{-\lambda_{2n}x} = \left(\sqrt{\left| \frac{1-\alpha}{1+\alpha} \right|} \right)^x e^{in\alpha\pi x} \\ + \left(\sqrt{\left| \frac{1+\alpha}{1-\alpha} \right|} \right)^x e^{-in\alpha\pi x} + o\left(\frac{1}{|n|}\right) \\ \lambda_{2n}^{-1} p'_{2n}(x) = \left(\sqrt{\left| \frac{1-\alpha}{1+\alpha} \right|} \right)^x e^{in\alpha\pi x} \\ - \left(\sqrt{\left| \frac{1+\alpha}{1-\alpha} \right|} \right)^x e^{-in\alpha\pi x} + o\left(\frac{1}{|n|}\right), \quad |n| \rightarrow \infty \end{cases} \quad (3.7)$$

$$\begin{cases} F_n(x) = (\lambda_{1n}^{-1} f_{1n}(x), f_{1n}(x)), \quad n \in \mathbb{Z} \\ f_{1n}(x) = e^{\lambda_{1n}x} - e^{-\lambda_{1n}x} = \left(\sqrt{\left| \frac{1-\alpha}{1+\alpha} \right|} \right)^x e^{in\alpha\pi x} \\ - \left(\sqrt{\left| \frac{1+\alpha}{1-\alpha} \right|} \right)^x e^{-in\alpha\pi x} \\ \lambda_{1n}^{-1} f'_{1n}(x) = \left(\sqrt{\left| \frac{1-\alpha}{1+\alpha} \right|} \right)^x e^{in\alpha\pi x} \\ + \left(\sqrt{\left| \frac{1+\alpha}{1-\alpha} \right|} \right)^x e^{-in\alpha\pi x} \end{cases} \quad (3.8)$$

$$\begin{cases} \tilde{F}_n(x) = (\lambda_{2n}^{-1} \phi_{2n}(x), \phi_{2n}(x)) \\ \phi_{2n}(x) = e^{\lambda_{2n}x} - e^{-\lambda_{2n}x} = \left(\sqrt{\left| \frac{1-\alpha}{1+\alpha} \right|} \right)^x e^{in\alpha\pi x} \\ - \left(\sqrt{\left| \frac{1+\alpha}{1-\alpha} \right|} \right)^x e^{-in\alpha\pi x} + o\left(\frac{1}{|n|}\right) \\ \lambda_{2n}^{-1} \phi'_{2n}(x) = \left(\sqrt{\left| \frac{1-\alpha}{1+\alpha} \right|} \right)^x e^{in\alpha\pi x} \\ + \left(\sqrt{\left| \frac{1+\alpha}{1-\alpha} \right|} \right)^x e^{-in\alpha\pi x} + o\left(\frac{1}{|n|}\right), \quad |n| \rightarrow \infty. \end{cases} \quad (3.9)$$

Proof: Suppose $\mathbb{A}(f, g, \phi, \psi) = \lambda(f, g, \phi, \psi) \neq 0$. Then, we have $g = \lambda f$, $\psi = \lambda \phi$ and (f, ϕ) satisfies

$$\begin{cases} f''(x) = \lambda^2 f(x) \\ f'(0) = \alpha \lambda f(0) + \beta f(0) + \phi'(0), f'(1) = -\alpha \lambda f(1) \\ \phi'' = \lambda^2 \phi \\ \phi(0) = 0, \phi'(1) = -\alpha \lambda f(1). \end{cases} \quad (3.10)$$

Let $p(x) = f(1-x) - \phi(1-x)$. Then, p satisfies

$$\begin{cases} p''(x) = \lambda^2 p(x) \\ p'(0) = 0, p'(1) = -\alpha \lambda p(1) - \beta p(1). \end{cases} \quad (3.11)$$

Now, if $p \equiv 0$, that is, $f \equiv \phi$. Then, $f \neq 0$ satisfies

$$\begin{cases} f''(x) = \lambda^2 f(x) \\ f(0) = 0, f'(1) = -\alpha \lambda f(1). \end{cases} \quad (3.12)$$

We get the eigenvalues

$$\lambda_{1n} = \frac{1}{2} \log \left| \frac{1-\alpha}{1+\alpha} \right| + n_\alpha \pi i, \quad \text{Re } \lambda_{1n} < 0, \quad n \in \mathbb{Z}.$$

This is the first expression of (3.3). The corresponding eigenvectors are

$$W_{1n}(x) = (\lambda_{1n}^{-1} f_{1n}(x), f_{1n}(x), \lambda_{1n}^{-1} f_{1n}(x), f_{1n}(x)) \quad n \in \mathbb{Z}$$

where f_{1n} is given by (3.8). This is the first expression of (3.5).

When $p \neq 0$, multiply by \bar{p} on both sides of the first equation of (3.11) and integrate by parts to give

$$\operatorname{Re} \lambda < 0. \quad (3.13)$$

The characteristic equation of (3.11) is

$$\begin{aligned} e^{2\lambda} &= \frac{(1-\alpha)\lambda - \beta}{(1+\alpha)\lambda + \beta} \\ &= \frac{1-\alpha}{1+\alpha} - \frac{2\beta}{(1+\alpha)^2\lambda} + o(|\lambda|^{-2}), \quad |\lambda| \rightarrow \infty. \end{aligned}$$

Solve the previous equation by virtue of Rouché's theorem to give the second expression of (3.3)

$$\lambda_{2n} = \frac{1}{2} \log \left| \frac{1-\alpha}{1+\alpha} \right| + n^\alpha \pi i + o\left(\frac{1}{|n|}\right), \quad |n| \rightarrow \infty.$$

The corresponding solutions p_{2n} to (3.11) are expressed by (3.7), which holds for $x \in [0, 1]$ uniformly. It follows from (3.10) that the corresponding eigenfunctions are

$$W_{2n}(x) = (\lambda_{2n}^{-1} f_{2n}(x), f_{2n}(x), \lambda_{2n}^{-1} \phi_{2n}(x), \phi_{2n}(x))$$

where $\phi_{2n}(x)$ is given in (3.9). A simple computation shows that

$$f_{2n}(x) = \phi_{2n}(x) + c_{2n} p_{2n}(1-x)$$

where c_{2n} is given in (3.6). The second expression of (3.5) is derived. The proof is complete. \square

Let us recall that $W \in D(\mathbb{A})$ is said to be a generalized eigenfunction of \mathbb{A} associated with the eigenvalue λ if there is an integer $\ell \geq 1$ such that $(\lambda - \mathbb{A})^\ell W = 0$. The integer $m_{(a)}(\lambda) = \sup_{\ell \in \mathbb{N}} \dim \ker(\lambda - \mathbb{A})^\ell$ is called the algebraic multiplicity of λ . An eigenvalue λ is said algebraically simple if $m_{(a)}(\lambda) = 1$.

The sequence $\{W_n\}_{n \in \mathbb{Z}}$ is called a basis for X if to each element $W \in X$ corresponds a unique sequence of numbers $\{c_n\}$ such that the series

$$W = \sum_{n \in \mathbb{Z}} c_n W_n \quad (3.14)$$

is convergent with respect to the norm of X . $\{W_n\}_{n \in \mathbb{Z}}$ is called a Riesz basis for X if

- $\overline{\operatorname{span}}\{W_n\} = X$;
- there exist some positive constants m_1 and m_2 such that for any numbers $c_n, n \in \mathcal{I}$, where \mathcal{I} is any finite subset of \mathbb{Z} , the following holds:

$$m_1 \sum_{n \in \mathcal{I}} |c_n|^2 \leq \left\| \sum_{n \in \mathcal{I}} c_n W_n \right\|^2 \leq m_2 \sum_{n \in \mathcal{I}} |c_n|^2.$$

We refer to [26] for more details on Riesz basis.

Remark 1: In designing the output feedback control law, we use the same gain α as that in (2.1). Actually, we can design $u(t) = -\gamma \hat{w}_t(1, t)$ for any $\gamma < 0$. In this case, by the same spectral analysis as in the proof of Lemma 2, the spectrum are still as that in (3.3) but λ_{1n} is replaced by $\lambda_{1n} = (1/2) \log |(1-\gamma)/(1+\gamma)| + n_\gamma \pi i$. So the decay rate is still dominated by λ_{2n} . It does not bring any advantage for the stabilization.

Theorem 3: Suppose $\alpha \neq 1$. Then, the following assertions hold true:

- \mathbb{A} has a family of generalized eigenfunctions, which forms a Riesz basis for X ;
- all eigenvalues of \mathbb{A} with sufficiently large module are algebraic simple;
- \mathbb{A} generates a C_0 -semigroup on X ;
- the spectrum-determined growth condition holds for $e^{\mathbb{A}t}$: $\omega(\mathbb{A}) = S(\mathbb{A}) < 0$, where $\omega(\mathbb{A})$ is the growth order of $e^{\mathbb{A}t}$;
- the semigroup $e^{\mathbb{A}t}$ is exponentially stable.

Proof: The relationship between iii) and iv) and i) and ii) is well-known, which can be found for the explanations of (5.1) of [22]). Hence ii)–v) are consequences of (i) and (3.3), (3.13) and the following arguments for the proof of assertion (i). We consider here only $0 < \alpha < 1$ since the case of $\alpha > 1$ can be treated similarly.

It is known that there exists a set of generalized eigenfunctions for the operator \mathbf{A} defined by (2.2), which forms a Riesz basis for \mathbf{H} (it can be shown simply by the asymptotic expression obtained in Lemma 2 and the abstract result in [7]). Its characteristic equation is just (3.11) and all generalized eigenfunctions are denoted by $\{G_n(x)\}_{n \in \mathbb{Z}_0}$, where \mathbb{Z}_0 is a set of index with the property that there exists a large integer N such that $\mathbb{Z}_0 - \{n \mid |n| > N\}$ is a finite set. The asymptotic eigenfunctions for $G_n(x)$ are given by (3.6). From the arguments in the sequel, we may assume without loss of generality that all $W_{2n}(x)$ have the same form of (3.5): $\{W_{2n}(x) = (\tilde{F}_n(x) + G_n(x), \tilde{F}_n(x))\}_{n \in \mathbb{Z}_0}$. All eigenvalues of \mathbf{A} with large modulus are algebraically simple.

Let us define another operator $\tilde{\mathbf{A}}$ in H

$$\begin{cases} \tilde{\mathbf{A}}(f, g) = (g, f'') \\ D(\tilde{\mathbf{A}}) = \{(f, g) \in H \mid \tilde{\mathbf{A}}(f, g) \in H, f'(1) = -\alpha g(1)\}. \end{cases}$$

Its characteristic equation is (3.12). It is well known that the eigenfunctions $\{F_n(x)\}_{n \in \mathbb{Z}}$ of $\tilde{\mathbf{A}}$ given by (3.8) form a Riesz basis for H [17].

Now, define an isometric isomorphism $\mathbb{T}_1 : X \rightarrow (L^2(0, 1))^2 \times \mathcal{C} \times (L^2(0, 1))^2$ by

$$\mathbb{T}_1(f, g, \phi, \psi) = (f', g, \sqrt{\beta} f(0), \phi', \psi) \quad \forall (f, g, \phi, \psi) \in X.$$

Then, the proof will be accomplished if we can show that

$$\left\{ \left\{ \mathbb{T}_1 \begin{pmatrix} F_n(x) \\ F_n(x) \end{pmatrix} \right\}_{n \in \mathbb{Z}}, \left\{ \mathbb{T}_1 \begin{pmatrix} \tilde{F}_n(x) + G_n(x) \\ \tilde{F}_n(x) \end{pmatrix} \right\}_{n \in \mathbb{Z}_0} \right\}$$

forms a Riesz basis for $(L^2(0, 1))^2 \times \mathcal{C} \times (L^2(0, 1))^2$.

Next, define a bounded invertible transformation \mathbb{T}_0 on $L^2(0, 1) \times L^2(0, 1)$:

$$\begin{aligned} \mathbb{T}_0(f, g)(x) &= (-g(1-x), f(1-x)) \\ &\quad \forall (f, g) \in L^2(0, 1) \times L^2(0, 1). \end{aligned}$$

Then, it is easy to see that (we use the column and row notations alternatively without confusion)

$$\begin{aligned} & \begin{pmatrix} I & 0 & -I \\ 0 & 1 & 0 \\ c_{2n}^{-1}\mathbb{T}_0 & 0 & I - c_{2n}^{-1}\mathbb{T}_0 \end{pmatrix} \mathbb{T}_1 \begin{pmatrix} F_n(x) \\ F_n(x) \end{pmatrix} \\ &= \mathbb{T}_1 \begin{pmatrix} 0 \\ F_n(x) \end{pmatrix} \\ & \begin{pmatrix} I & 0 & -I \\ 0 & 1 & 0 \\ c_{2n}^{-1}\mathbb{T}_0 & 0 & I - c_{2n}^{-1}\mathbb{T}_0 \end{pmatrix} \cdot \\ & \mathbb{T}_1 \begin{pmatrix} \tilde{F}_n(x) + G_n(x) \\ \tilde{F}_n(x) \end{pmatrix} \\ &= \mathbb{T}_1 \begin{pmatrix} G_n(x) \\ 0 \end{pmatrix} + o\left(\frac{1}{|n|}\right) \end{aligned}$$

where c_{2n} is defined in (3.6) and I is the identity on $(L^2(0,1))^2$. It is obvious that $\{(0, F_n(x))\}_{n \in \mathbb{Z}}, \{(G_n(x), 0)\}_{n \in \mathbb{Z}_0}$ forms a Riesz basis for X and so

$$\left\{ \left\{ \mathbb{T}_1 \begin{pmatrix} 0 \\ F_n(x) \end{pmatrix} \right\}_{n \in \mathbb{Z}}, \left\{ \mathbb{T}_1 \begin{pmatrix} G_n(x) \\ 0 \end{pmatrix} \right\}_{n \in \mathbb{Z}_0} \right\}$$

forms a Riesz basis for $(L^2(0,1))^2 \times \mathcal{C} \times (L^2(0,1))^2$. By classical Bari's theorem [11] and (3.15), it follows that

$$\left\{ \left\{ \begin{pmatrix} I & 0 & -I \\ 0 & 1 & 0 \\ c_{2n}^{-1}\mathbb{T}_0 & 0 & I - c_{2n}^{-1}\mathbb{T}_0 \end{pmatrix} \mathbb{T}_1 \begin{pmatrix} F_n(x) \\ F_n(x) \end{pmatrix} \right\}_{n \in \mathbb{Z}}, \left\{ \begin{pmatrix} I & 0 & -I \\ 0 & 1 & 0 \\ c_{2n}^{-1}\mathbb{T}_0 & 0 & I - c_{2n}^{-1}\mathbb{T}_0 \end{pmatrix} \mathbb{T}_1 \right. \right. \\ \left. \left. \times \begin{pmatrix} \tilde{F}_n(x) + G_n(x) \\ \tilde{F}_n(x) \end{pmatrix} \right\}_{n \in \mathbb{Z}_0} \right\}$$

forms a Riesz basis for $(L^2(0,1))^2 \times \mathcal{C} \times (L^2(0,1))^2$. Notice that for all $n \in \mathbb{Z}$ or $n \in \mathbb{Z}_0$, the transformations

$$\begin{pmatrix} I & 0 & -I \\ 0 & 1 & 0 \\ c_{2n}^{-1}\mathbb{T}_0 & 0 & I - c_{2n}^{-1}\mathbb{T}_0 \end{pmatrix}$$

and their inverses are uniformly bounded in $(L^2(0,1))^2 \times \mathcal{C} \times (L^2(0,1))^2$ with respect to n . Therefore

$$\left\{ \left\{ \mathbb{T}_1 \begin{pmatrix} F_n(x) \\ F_n(x) \end{pmatrix} \right\}_{n \in \mathbb{Z}}, \left\{ \mathbb{T}_1 \begin{pmatrix} \tilde{F}_n(x) + G_n(x) \\ \tilde{F}_n(x) \end{pmatrix} \right\}_{n \in \mathbb{Z}_0} \right\}$$

forms a Riesz basis for $(L^2(0,1))^2 \times \mathcal{C} \times (L^2(0,1))^2$. The proof is complete. \square

Finally, we indicate that Theorem 3 should be still true even for $\alpha = 1$. There is some other method to treat the exponential stability for nondissipative PDEs. We refer to [9] for the progress on this aspect. In order to understand the case of $\alpha = 1$ intuitively, we use an Legendre spectral method to approximate (3.11) and obtain two figures. Fig. 1 demonstrates the functional relation of $S(\mathcal{A})$ with respect to β where we take $\alpha = 0.7$. It recommends the choice of small β . Fig. 2

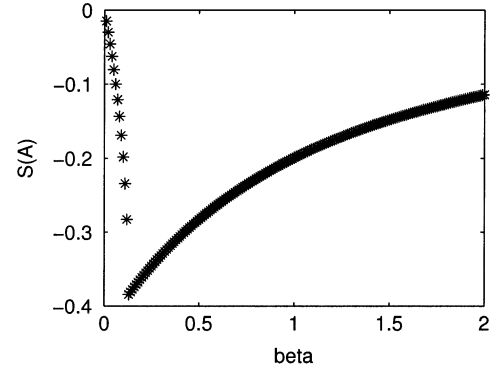


Fig. 1. Functional relation of $S(\mathcal{A})$ with respect to β with $\alpha = 0.7$.

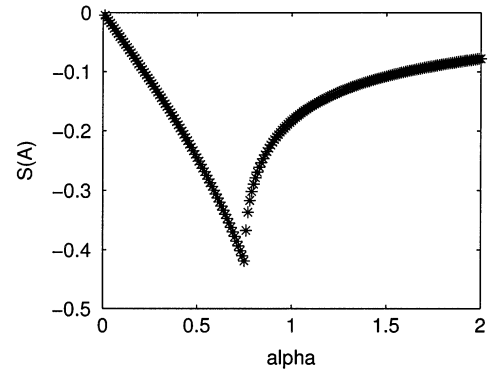


Fig. 2. Functional relation of $S(\mathcal{A})$ with respect to α with $\beta = 0.15$.

demonstrates the functional relation of $S(\mathcal{A})$ with respect to α where we take $\beta = 0.15$, which recommends the choice of α around one.

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